

FREQUENCY – DEPENDENT DIELECTRIC PROPERTIES OF $TiGa_{1-x}Co_xS_2$ SINGLE CRYSTALS

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Laylı $TiGa_{1-x}Co_xS_2$ ($x = 0,005; 0,01$) monokristallarda dielektrik xarakteristikaların tezlikdən asılılıqları və $f = 5 \cdot 10^4 \div 3,5 \cdot 10^7$ Hz tezlik diapazonunda laylara perpendikulyar istiqamətdə ac-keçiriciliyi (σ_{ac}) tədqiq edilmişdir. Müəyyən edilmişdir ki, $TiGa_{1-x}Co_xS_2$ kristallarda elektrik keçiriciliyində olan itkilərlə yanaşı həmçinin relaksasiya ilə bağlı olan itkilərdə mövcuddur.

Müəyyən edilmişdir ki, $5 \cdot 10^4 \div 2 \cdot 10^7$ Hz tezlik sahəsində $TiGa_{1-x}Co_xS_2$ kristallarda yükün köçürülməsi, Fermi səviyyəsi yaxınlığında lokalizə olunmuş hallar üzrə, sıçrayış mexanizmi ilə baş verir. Bu səviyyələrin sıxlığı (N_F), energetik səpələnməsi (ΔE), sıçrayışların orta vaxtı (τ) və məsafəsi (R) qiymətləndirilmişdir. $TiGa_{1-x}Co_xS_2$ monokristalların tərkibinin onların dielektrik əmsallarına təsiri təyin edilmişdir. Müəyyən olunmuşdur ki, monokristallarda kobaltın miqdarı artdıqca N_F və ΔE yüksəlir, τ və R isə azalır.

В слоистых монокристаллах $TiGa_{1-x}Co_xS_2$ ($x = 0,005; 0,01$) изучены частотные зависимости диэлектрических характеристик и ac-проводимости (σ_{ac}) поперек слоёв в частотном диапазоне $f = 5 \cdot 10^4 \div 3,5 \cdot 10^7$ Гц. Установлено, что в кристаллах $TiGa_{1-x}Co_xS_2$ наряду с потерями на электропроводность имеют место также релаксационные потери. Обнаружено, что в частотной области $5 \cdot 10^4 \div 2 \cdot 10^7$ Гц в кристаллах $TiGa_{1-x}Co_xS_2$ имеет место прыжковый механизм переноса заряда по локализованным вблизи уровня Ферми состояниям. Оценены: плотность (N_F) и энергетический разброс (ΔE) этих состояний, среднее время (τ) и расстояние (R) прыжков. Изучено влияние состава монокристаллов $TiGa_{1-x}Co_xS_2$ на их диэлектрические коэффициенты. Установлено, что с ростом содержания кобальта N_F и ΔE возрастают, а τ и R уменьшаются.

Frequency dependence of the dissipation factor $\tan\delta$, the permittivity ϵ , and the ac conductivity σ_{ac} across the layers in the frequency range $f = 5 \cdot 10^4 \div 3,5 \cdot 10^7$ Hz were studied in layered $TiGa_{1-x}Co_xS_2$ single crystals ($x = 0,005; 0,01$). The results demonstrate that dielectric dispersion in the $TiGa_{1-x}Co_xS_2$ crystals has a relaxation nature. It was established that the mechanism of ac charge transport across the layers in $TiGa_{1-x}Co_xS_2$ single crystals in the frequency range from $5 \cdot 10^4$ to $2 \cdot 10^7$ Hz is hopping over localized states near the Fermi level. The Fermi level density of states (N_F), the spread of their energies (ΔE), and the mean hop distance (R) and time (τ) have been estimated. The present results demonstrate that partial substitution of cobalt for gallium offers the possibility of tuning dielectric properties of obtained crystals.

Layered $TiGaS_2$ single crystals characterized by sufficiently high density of localized states at the Fermi level (N_F). The N_F value calculated from the experimental results of dc- and ac- conductivity measurements of $TiGaS_2$ single crystals is $\sim 2 \cdot 10^{18} \text{ eV}^{-1} \cdot \text{cm}^{-3}$ [1,2]. Doping of these crystals allows one to tune their physical properties. For example, in [3] it was shown that partial manganese substitution for gallium in $TiGaS_2$ leads to significant decrease of dielectric permittivity, mean hop distance and time, and also to increase of density of localized states near Fermi level.

The purpose of this work was to study the effect of partial cobalt substitution for gallium on the ac dielectric properties of $TiGaS_2$ single crystals.

Interaction in $TiGaS_2 - TiCoS_2$ system has been investigated by the methods of differential thermal and roentgenography analyses in our previous work [4]. Synthesis of initial $TiGaS_2$ and $TiCoS_2$ compounds produced in ampoules, evacuated to the pressure 10^{-3} Pa. Ampoules were made of melted quartz tubes. Synthesis of samples has been carried out at interactions of initial elements (Ti, Ga, Co, S) of high purity degree ($> 99,999\%$). $TiGa_{1-x}Co_xS_2$ single crystals were growth by Bridgman method.

Measurements of the dielectric properties of $TiGa_{1-x}Co_xS_2$ ($x = 0,005; 0,01$) single crystals were performed at fixed frequencies in the range of $5 \cdot 10^4 \div 3,5 \cdot 10^7$ Hz by the resonant method using a TESLA BM560 Q-meter.

The single crystal $TiGa_{1-x}Co_xS_2$ samples for dielectric measurements had the form of planar capacitors normal to the C axis of the crystals, with silver-paste electrodes. The thickness of the single crystal $TiGa_{1-x}Co_xS_2$ samples was $(1,0 \div 1,4) \cdot 10^{-2}$ cm and the area of the capacitor plates was $7,6 \cdot 10^2$

cm^2 . All dielectric measurements were performed at $T = 300\text{K}$. The accuracy in determining the resonance capacitance and the quality factor $Q = 1/\tan\delta$ of the measuring circuit was limited by errors related to the resolution of the device readings. The accuracy of the capacitor graduation was $\pm 0,1$ pF. The reproducibility of the resonance position was $\pm 0,2$ pF in capacitance and $\pm(1,0 \div 1,5)$ scale divisions in quality factor. The largest deviations were 3 - 4% in ϵ and 7% in $\tan\delta$.

Fig. 1 shows the experimental frequency dependences of the dissipation factor $\tan\delta$ for $TiGa_{1-x}Co_xS_2$ ($x = 0; 0,005; 0,01$) single crystals at $5 \cdot 10^4 \div 3,5 \cdot 10^7$ Hz.

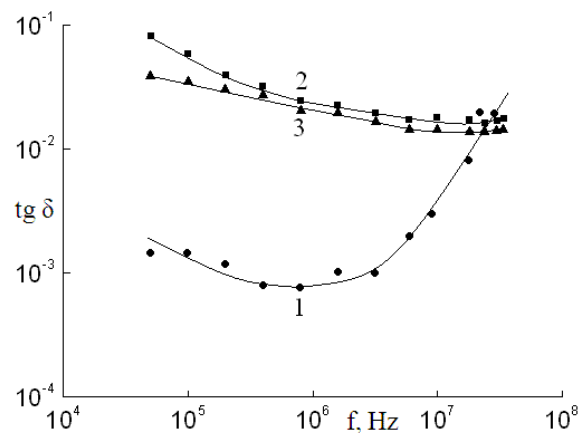


Fig. 1. Frequency dependences of the dissipation factor of $TiGa_{1-x}Co_xS_2$ single crystals: 1 – $x = 0$; 2 – $x = 0,005$; 3 – $x = 0,01$ ($T = 300\text{K}$).

The $\tan\delta(f)$ curves have a steadily descending branch. This experimental result is evidence of the fact that conductivity loss becomes the main dielectric loss mechanism in studied $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ single crystals. In TlGaS_2 at $f > 10^6$ Hz significant increase of $\tan\delta$ takes place up to $3 \cdot 10^7$ Hz. This experimental result points to relaxation losses. In $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ ($x = 0.005; 0.01$) single crystals insignificant increase of $\tan\delta$ takes place at $f > 2 \cdot 10^7$ Hz.

We also measured the electric capacitance of $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ samples in the frequency range $5 \cdot 10^4 \div 3.5 \cdot 10^7$ Hz. Using the measured capacities of these samples, we calculated the permittivity ϵ at different frequencies ($\epsilon = \text{CL}/\epsilon_0\text{S}$; L is thickness of crystal; S is contact area; ϵ_0 is dielectric constant).

Fig. 2 shows the frequency dependences of 300 - K dielectric permittivity for $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ single crystals with $x = 0; 0.005$, and 0.01 .

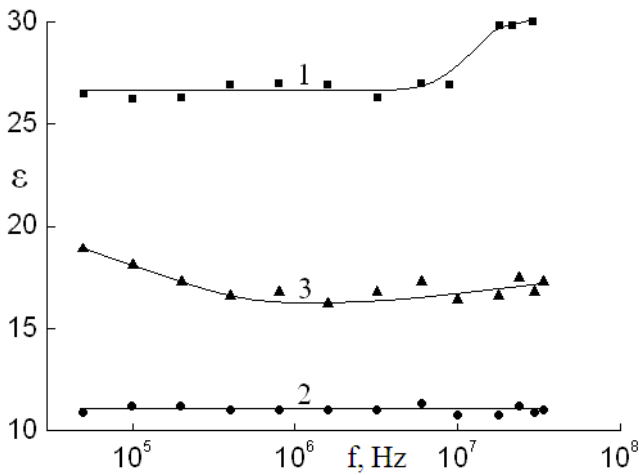


Fig. 2. Frequency dispersion of the dielectric permittivity of TlGaS_2 (1); $\text{TlGa}_{0.995}\text{Co}_{0.005}\text{S}_2$ (2), and $\text{TlGa}_{0.99}\text{Co}_{0.01}\text{S}_2$ (3).

Partial Ga^{3+} substitution for cobalt in TlGaS_2 crystals gives to considerable decreasing of ϵ at all frequencies. The ϵ values of TlGaS_2 and $\text{TlGa}_{0.995}\text{Co}_{0.005}\text{S}_2$ are seen to vary from 26 to 30 and from 10.7 to 11.3 respectively over the entire frequency range studied, with no significant dispersion (Fig. 2, curves 1 and 2). But in $\text{TlGa}_{0.99}\text{Co}_{0.01}\text{S}_2$ single crystal there was observed the monotonic reduction in the dielectric permittivity with increasing frequency (Fig.2, curve 3). This experimental result attests to the relaxation nature of the dispersion.

Fig. 3 shows the experimentally measured frequency dependence of the ac conductivity of $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ single crystals at $T=300\text{K}$. The σ_{ac} of TlGaS_2 (Fig. 3, curve 1) varies as $f^{0.8}$ in the frequency range $5 \cdot 10^4$ to 10^6 Hz and as f^2 in the range 10^6 to $3 \cdot 10^7$ Hz. As is known [5], conductivity proportional to f^2 is due to optical transitions in semiconductors and prevails at high frequencies.

The σ_{ac} of the $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ ($x = 0.005; 0.01$) samples (Fig. 3, curves 2 and 3) varies as $f^{0.8}$ at $f = 5 \cdot 10^4 \div 10^7$ and $5 \cdot 10^4 \div 2 \cdot 10^7$ Hz, respectively. At more high frequencies $\sigma_{ac}(f)$ – dependence of these crystals was super linear.

The $\sigma_{ac} \sim f^{0.8}$ dependence indicates that the mechanism of charge transport is hopping over localized states near Fermi

level [5]. This charge transport mechanism is characterized by the following expression obtained in [6]:

$$\sigma_{ac}(f) = \frac{\pi^3}{96} e^2 k T N_F^2 a^5 f \left[\ln\left(\frac{v_{ph}}{f}\right) \right]^4 \quad (1)$$

where e is the elementary charge, k is the Boltzmann constant, N_F is the density of localized states near the Fermi level, a is the localization length, and v_{ph} is the phonon frequency. Using expression (1), we can calculate the density of states at the Fermi level from the measured values of the conductivity. N_F was determined to be $7.9 \cdot 10^{18}$ and $9.1 \cdot 10^{18} \text{ eV}^{-1} \cdot \text{cm}^{-3}$ in $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ single crystals with $x = 0.005$, and 0.01 , respectively. As it was shown earlier, in TlGaS_2 N_F was determined to be $2.1 \cdot 10^{18} \text{ eV}^{-1} \cdot \text{cm}^{-3}$ [2]. In calculating N_F , the localization radius in $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ was taken to be $a = 14\text{\AA}$, by analogy with gallium sulfide [7], a binary analog of TlGaS_2 . The value of v_{ph} in TlGaS_2 is approximately 10^{12} Hz [8].

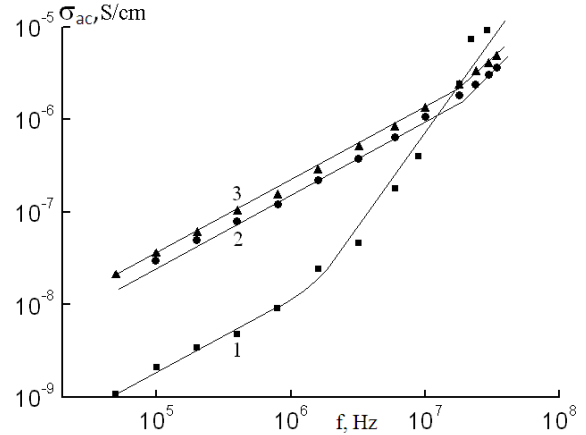


Fig. 3. Conductivity as a function of frequency for TlGaS_2 (1), $\text{TlGa}_{0.995}\text{Co}_{0.005}\text{S}_2$ (2), and $\text{TlGa}_{0.99}\text{Co}_{0.01}\text{S}_2$ (3).

According to the theory of hopping conduction [5], the mean hop distance in an applied ac electric field is given by

$$R = (1/2\alpha) \ln(v_{ph}/f), \quad (2)$$

where $1/f = \tau$ is the mean hop time, $\alpha = 1/a$ is the decay parameter of the wave function of a localized charge carrier, $\psi \sim e^{-\alpha r}$.

The R values calculated by Eq. (2) for $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ single crystals with $x=0.005$ and 0.01 are equal to 86 and 81\AA , respectively. For TlGaS_2 , $R = 103\text{\AA}$ [2]. In $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ crystals, R is ~ 6 times greater than the average distance between localization centers (in TlGaS_2 , $R/a = 7.3$ [2]). The mean hop time in $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ crystals is $0.2\mu\text{s}$ for $x = 0.005$ and $0.1\mu\text{s}$ for $x = 0.01$ (in TlGaS_2 , $\tau = 2.0\mu\text{s}$ [2]).

Using the relation [5]:

$$\Delta E = 3 / 2\pi R^3 N_F \quad (3)$$

we evaluate the spread of the Fermi-level states in energy. We obtained for $\text{TlGa}_{1-x}\text{Co}_x\text{S}_2$ $\Delta E = 9.0 \cdot 10^{-2} \text{ eV}$ at $x = 0.005$ and $9.7 \cdot 10^{-2} \text{ eV}$ at $x = 0.01$.

The concentration of deep traps determining the ac conductivity of TlGaS_2 ,

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$$N_t = N_F \cdot \Delta E \quad (4)$$

is $7.1 \cdot 10^{17} \text{cm}^{-3}$ at $x = 0.005$ and $8.8 \cdot 10^{17} \text{cm}^{-3}$ at $x = 0.01$.

Parameters of TlGa_{1-x}Co_xS₂ single crystals obtained from high-frequency dielectric measurements are given in Table:

Table

Crystal composition	$N_F, \text{eV}^{-1} \cdot \text{cm}^{-3}$	$R, \text{\AA}$	τ, s	R/a
TlGaS ₂	$2.1 \cdot 10^{18}$	103	2.0	7.3
TlGa _{0.995} Co _{0.005} S ₂	$7.9 \cdot 10^{18}$	86	0.2	6.1

TlGa _{0.99} Co _{0.01} S ₂	$9.1 \cdot 10^{18}$	81	0.1	5.8
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CONCLUSIONS

The dielectric properties of TlGa_{1-x}Co_xS₂ single crystals were studied in a broad frequency range. The nature of dielectric losses and the mechanism of charge transport in the crystals was established, and the Fermi-level density of states, the spread of their energies, the mean hop distance and time, and the concentration of deep traps determining the ac conductivity of the crystals were estimated.

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| <p>[1]. <i>S.N. Mustafaeva, V.A. Aliev and M.M. Asadov</i> DC hopping conduction in TlGaS₂ and TlInS₂ single crystals // <i>Fiz. Tverd. Tela (S. - Peterburg)</i>, 1998, v. 40, No 4, pp. 612– 615.</p> <p>[2]. [<i>S.N. Mustafaeva</i> Dielectric dispersion in TlGaS₂ layered crystals // <i>Fiz. Tverd. Tela (S. - Peterburg)</i>, 2004, v. 46, No 6, pp. 979 – 981.</p> <p>[3]. [<i>S.N. Mustafaeva</i> Dielectric properties of TlGa_{1-x}Mn_xS₂ ($0 \leq x \leq 0.03$) single crystals // <i>Inorganic Materials</i>, 2006, v. 42, No 5, pp. 470 – 473.</p> <p>[4]. <i>S.N. Mustafaeva, E.M. Kerimova, F.M. Seidov, A.I. Jabbarly</i> Novel thermoelectric materials on the base of TlGaS₂ – TlCoS₂ // <i>Abstracts of 13th International Conference on Ternary and Multinary</i></p> | <p><i>Compounds (ICTMC-13)</i>. Paris. France. Oct. 14 – 18, 2002. p. 125.</p> <p>[5]. <i>N.F. Mott and E.A. Devis</i> <i>Electronic Processes in Non-Crystalline Materials</i>. / Oxford: Clarendon, 1971.</p> <p>[6]. <i>M. Pollak</i> Frequency dependence of conductivity in amorphous solids // <i>Philos. Mag.</i>, 1971, v. 23, pp. 519 – 542</p> <p>[7]. <i>V. Augelli, C. Manfredotti, R. Murri et al.</i> Anomalous impurity conductivity in n-GaSe and n-GaS // <i>Nuovo Cimento Soc. Ital. Fis., B</i>, 1977, v. 38, pp. 327 – 336.</p> <p>[8]. <i>K.R. Allakhverdiev, E.A. Vinogradov, R. Kh. Nani et al.</i> Vibrational spectra of TlGaS₂, TlGaSe₂, and β-TlInS₂ crystals. / <i>Fizicheskie svoistva slozhnykh poluprovodnikov (Physical properties of Compound semiconductors)</i>. Baku: Elm, 1982, pp. 55 – 63.</p> |
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Daxil olunub: 01.07.2007