# OPTIMIZATION OF PARAMETERS AND THE GEOMETRICAL SIZES THE PRECISION STABILIZER OF THE ALTERNATING CURRENT 

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#### Abstract

In work optimization of parameters and the geometrical sizes of the precision stabilizer of an alternating current is considered. The criterion of optimization chooses optimum values of density of currents in windings. On the basis of the received optimum parities recommendations at the choice of parameters of the stabilizer are developed.


Keywords: stabilizer, optimization, magnetic fields, galvanic baths.

## I. INTRODUCTION

Test beds, galvanic baths and installations where scientific researches on the stabilized magnetic fields are carried out require precision stabilizers of an alternating current. In the present work optimization of parameters and the geometrical sizes of the stabilizer of an alternating current with short-circuited levitation a winding (LW) [1] is considered. In the given work as criterion of optimization preset values of density of currents in windings are accepted.

## II. MAIN PART

According to work [2] for a preset value of density of current $\mathrm{J}_{2}$ in LW it is possible to write down:

$$
\begin{equation*}
\tau_{\mathrm{T} 2}=\mathrm{A}_{\mathrm{T} 2} \mathrm{~J}_{2}\left(\frac{\mu_{\mathrm{o}} 1_{\mathrm{c} 2}}{\lambda \mathrm{~S}_{\mathrm{a} \hat{\mathrm{a}} \hat{\mathrm{e}} 2}^{2}}\right)=\mathrm{J}_{2} \mathrm{~m}_{12}^{*} \mathrm{~A}_{\mathrm{T} 2} \mathrm{~A}_{\mathrm{C} 2} \sqrt{\frac{\mathrm{C}}{\mathrm{n}_{\mathrm{e} 2}}}, \tag{1}
\end{equation*}
$$

Where are designated:

$$
\begin{aligned}
& \mathrm{A}_{\mathrm{T} 2}=\frac{\mathrm{b}_{2} \rho_{20}}{\mathrm{~K}_{\mathrm{T}}} \sqrt{\frac{2 \mathrm{~g} \gamma_{2} \mathrm{~K}_{32}}{\mu_{\mathrm{o}}}} ; \mathrm{A}_{\mathrm{C} 2}=\sqrt{\frac{\sigma_{2}}{\sigma_{\hat{\mathrm{a}}}}} ; \frac{1_{\mathrm{cp} 2}}{\lambda}=\frac{\mathrm{C}}{\mu_{\mathrm{o}}} \mathrm{~A}_{\mathrm{C} 2} ; \\
& \sigma_{2}=\frac{2 \mathrm{~m}_{\mathrm{a}}+\mathrm{m}_{\mathrm{c}} \mathrm{n}_{\mathrm{o} 2}+\mathrm{m}_{\mathrm{a}} \mathrm{~m}_{\mathrm{c}} \mathrm{n}_{\mathrm{o} 2}}{\mathrm{~m}_{\mathrm{a}} \mathrm{~m}_{\mathrm{c}} \mathrm{n}_{\mathrm{o} 2}}
\end{aligned}
$$

If to accept: $\gamma_{2}=n_{k} \gamma_{M}=1.1 \cdot 8.9 \cdot 10^{3}=9.79 \cdot 10^{3} \mathrm{Kг} / \mathrm{M}^{3} ; \rho_{20}$ $=1.7 \cdot 10^{-8} \mathrm{Om} / \mathrm{m} ; \mathrm{g}=9.8 \mathrm{~m} / \mathrm{c}^{2} ; \mathrm{K}_{\mathrm{T}}=13 ; \mathrm{K}_{32}=0.6 ; \mu_{\mathrm{o}}$ $=1.256 \cdot 10^{-6}{ }^{\mathrm{H}} / \mathrm{M} ; \mathrm{b}_{2}=0.98$ we shall receive: $\mathrm{A}_{\mathrm{T} 2}=$ $389 \cdot 10^{-6} ; \mathrm{A}_{\mathrm{C} 2}=1.082 \div 1.283$. From (1) it is easy to define size:
$\tilde{\mathrm{N}}=\left(\frac{\tau_{\mathrm{T} 2}}{\mathrm{~A}_{\mathrm{T} 2}}\right)^{2} \frac{\mathrm{n}_{\mathrm{e} 2}}{\mathrm{~J}_{2}^{2} \mathrm{~A}_{\mathrm{c} 2}^{2} \mathrm{~m}_{12}^{* 2}}$
(2) Using (2) it is
easy to find other sizes of the core and windings if values $\mathrm{J}_{1}$ are known, $\mathrm{J}_{2}$. With this purpose in the beginning we shall consider as density of currents $\mathrm{J}_{1}, \mathrm{~J}_{2}$ depending on temperatures $\tau_{1}, \tau_{2}$ change.

We Shall present force of weight LW $\mathrm{P}_{\mathrm{B}}$ and temperature $\tau_{1}$ in a following kind:

$$
\begin{gathered}
\mathrm{P}_{\hat{\mathrm{a}}}=2\left(\mathrm{gyb}_{2}\right)^{2} \frac{1_{\mathrm{cp} 2}^{2}}{\lambda \mathrm{~J}_{2}^{2}} \\
\tau_{1}=\mathrm{A}_{1} \mathrm{~J}_{1} \sqrt{\frac{\mathrm{~S}_{\mathrm{ol} 1} \mathrm{c}_{\mathrm{cp} 1}^{2}}{\lambda \mathrm{~S}_{\mathrm{ox} 1}}\left(\Delta_{\mathrm{ok}}+\alpha_{\mathrm{M}} \tau_{1}\right)=\mathrm{J}_{1} \frac{\mathrm{~K}_{\mathrm{p}}}{\mathrm{~b}_{2}} \mathrm{~A}_{\mathrm{C} 1} \mathrm{~A}_{\mathrm{T} 2} \sqrt{\frac{\mathrm{C}}{\mathrm{n}_{\mathrm{el}}}}},
\end{gathered}
$$

where:

$$
\begin{aligned}
& \mathrm{A}_{1}=\frac{\mathrm{K}_{\mathrm{p}}}{\mathrm{~K}_{\mathrm{T}}} \rho_{10} \sqrt{2 \mathrm{~g} \gamma \mathrm{~K}_{31}} ; \mathrm{A}_{\mathrm{C} 1}=\sqrt{\frac{\sigma_{1}}{\sigma_{\hat{\mathrm{a}}}}} ; \mathrm{A}_{\mathrm{T} 1}=\mathrm{A}_{\mathrm{T} 2} \frac{\mathrm{~K}_{\mathrm{p}}}{\mathrm{~b}_{2}} ; \\
& \mathrm{I}_{1} \mathrm{~W}_{1}=\sqrt{2 \mathrm{~g} \mathrm{\gamma K}_{31} \mathrm{~S}_{\mathrm{ol}} \frac{1_{\mathrm{cp} 1}}{\lambda}}
\end{aligned}
$$

Then

$$
\frac{\tau_{\mathrm{T} 1}}{\tau_{\mathrm{T} 2}}=\left(\frac{\mathrm{J}_{1}}{\mathrm{~J}_{2}}\right)\left(\frac{\mathrm{K}_{\mathrm{p}}}{\mathrm{~b}_{2}}\right)\left(\frac{\mathrm{A}_{\mathrm{C} 1}}{\mathrm{~A}_{\mathrm{C} 2}}\right) \sqrt{\frac{\mathrm{n}_{\mathrm{e} 2}}{\mathrm{n}_{\mathrm{e} 1}}}
$$

As

$$
\mathrm{J}_{1} \mathrm{~K}_{\mathrm{p}}=\mathrm{J}_{1}+\mathrm{b}_{2} \mathrm{j}_{2} \mathrm{~K}_{\mathrm{c}} \mathrm{~K}_{\tau} ; \quad \frac{\mathrm{A}_{\mathrm{C} 1}}{\mathrm{~A}_{\mathrm{C} 2}}=\sqrt{\frac{\sigma_{1}}{\sigma_{2}}}=\frac{1}{\sqrt{\mathrm{~K}_{\mathrm{c}}}},
$$

That we write down:

$$
\frac{\tau_{\mathrm{T} 1}}{\tau_{\mathrm{T} 2}}=\left(\frac{\varepsilon_{12}}{\mathrm{~b}_{2}}+\mathrm{K}_{\mathrm{c}} \mathrm{~K}_{\tau}\right) \sqrt{\frac{\mathrm{n}_{\mathrm{e} 2}}{\mathrm{n}_{\mathrm{e} 1} \mathrm{~K}_{\mathrm{c}}}}, \quad\left(\text { where } \quad \varepsilon_{12}=\frac{\mathrm{J}_{1}}{\mathrm{~J}_{2}}\right)
$$

From here we shall receive:

$$
\varepsilon_{12}=\mathrm{b}_{2}\left(\frac{\tau_{\mathrm{T} 1}}{\tau_{\mathrm{T} 2}}\right) \sqrt{\mathrm{K}_{\mathrm{c}} \frac{\mathrm{n}_{\mathrm{e} 1}}{\mathrm{n}_{\mathrm{e} 2}}}-\mathrm{b}_{2} \mathrm{~K}_{\mathrm{c}} \mathrm{~K}_{\tau}
$$

Under condition of $\tau_{1}=\tau_{2}=\tau_{\mathrm{a}}$ for direct cores we shall receive:

$$
\left(\frac{\varepsilon_{12}}{\mathrm{~b}_{2}}+1\right) \sqrt{\frac{\mathrm{n}_{\mathrm{e} 2}}{\mathrm{n}_{\mathrm{e} 1}}}-1 ; \quad \text { If } \quad \hat{\mathrm{E}}_{\tilde{\mathrm{n}}}=\hat{\mathrm{E}}_{\tau}=1
$$

Then:
$\varepsilon_{12}=b_{2}\left(\sqrt{\frac{\mathrm{n}_{\mathrm{el}}}{\mathrm{n}_{\mathrm{e} 2}}}-1\right)=\mathrm{b}_{2}\left(\sqrt{\frac{\mathrm{~h}_{1}}{\mathrm{~h}_{2}}}-1\right) ; \quad$ where $\quad \frac{\mathrm{n}_{\mathrm{el}}}{\mathrm{n}_{\mathrm{e} 2}}=\frac{\mathrm{S}_{01}}{\mathrm{~S}_{\mathrm{o} 2}}=\frac{\mathrm{h}_{1}}{\mathrm{~h}_{2}}$
Thus, for direct cores under condition of $\tau_{1}=\tau_{2}=\tau_{\mathrm{a}}$ attitude $\mathrm{J}_{1} / \mathrm{J}_{2}$ depends only on the attitude of height of windings $h_{1} / h_{2}$, and always we take place $J_{1}<J_{2}$. On fig. 1 dependence $J_{1} / J_{2}=f\left(h_{1} / h_{2}\right)$ is resulted.
For step cores under condition of $\tau_{\mathrm{T} 1}=\tau_{\mathrm{T} 2}=\tau_{\mathrm{a}}$ we shall receive:

$$
\varepsilon_{12}=\mathrm{b}_{2}\left(\sqrt{\frac{\mathrm{n}_{\mathrm{e} 1}}{\mathrm{n}_{\mathrm{e} 2}} \mathrm{~K}_{\mathrm{c}}}-\mathrm{K}_{\mathrm{c}}\right)
$$

On figure. 2 and 3 dependences $\left|J_{1} / J_{2}\right|=f\left(n_{e 1} / n_{e 2}\right)$ are resulted, $\left|J_{1} / J_{2}\right|==f\left(K_{c}\right)$. On the basis of these dependences it is possible to draw following conclusions:

1. Because of performance of conditions of levitation $\mathrm{P}_{\mathrm{B}}=\mathrm{F}_{3}$ the density of a current in a motionless winding turns out less, than density of a current in LW.
2. For step cores $n_{e 1}<n_{e 2}$, therefore attitude $J_{1} / J_{2}$ at $n_{e 1} / n_{e 2}$ $=0.3-0.8$ makes $0.1-0.45$. With increase in factor of frequency rate LW: $\mathrm{n}_{\mathrm{e} 2}=\mathrm{h}_{2} / \mathrm{C}_{2}$ the density of a current in motionless winding $\mathrm{J}_{1}$ decreases.
With the purpose of definition of density of a current in LW from expressions:

$$
\mathrm{F}_{2}=\mathrm{K}_{32} \mathrm{n}_{\mathrm{e} 2} \mathrm{~J}_{2} \mathrm{C}_{2}^{2} ; \quad \mathrm{F}_{2}=\mathrm{b}_{2} \sqrt{\frac{2 \mathrm{P}_{\hat{\mathrm{a}}}}{\lambda}}=\mathrm{b}_{2} \mathrm{~F}_{1},
$$

We find thickness of winding:

$$
\mathrm{C}_{2}=\sqrt{\frac{\mathrm{b}_{2} \mathrm{~F}_{1}}{\mathrm{~K}_{32} \mathrm{n}_{\mathrm{e} 2} \mathrm{~J}_{2}}}=\mathrm{C}_{2}^{\prime}
$$



Fig.1. Dependence $\frac{\mathrm{J}_{1}}{\mathrm{~J}_{2}}=\mathrm{f}\left(\frac{\mathrm{h}_{1}}{\mathrm{~h}_{2}}\right)$


Fig. 2 Dependence at $\frac{J_{1}}{J_{2}}=f\left(\frac{n_{e 1}}{n_{e 2}}\right)$ various
values $K_{c}\left(K_{c}=0.9-1.0\right)$


Fig. 3 Dependence $\frac{J_{1}}{J_{2}}=f\left(K_{c}\right)$
This formula considers a condition of a levitation, but does not consider overheat LW. In view of an overheat, we shall receive:

$$
\begin{equation*}
\mathrm{C}_{2}=\frac{\mathrm{A}_{\mathrm{o} 2}}{\mathrm{~J}_{2}^{2}\left(\alpha_{\mathrm{M}}+\frac{\Delta_{\mathrm{ok}}}{\tau_{\text {ä } 2}}\right)}=\mathrm{C}_{2}^{\prime \prime}, \text { Here } \mathrm{A}_{\mathrm{o} 2}=\frac{\mathrm{K}_{\mathrm{T}}}{\mathrm{~K}_{32} \rho_{2 \mathrm{o}} \mathrm{~K}_{\mathrm{c} 2}} \tag{3}
\end{equation*}
$$

3. More effective way of increase in density of current $\mathrm{J}_{2}$ is reduction of a backlash With or thickness of winding $\mathrm{C}_{2}$. Thus factor $\mathrm{n}_{\mathrm{e} 2}=\mathrm{h}_{2} / \mathrm{C}_{2}$ and conductivity 1 increase.

$$
\tau_{\mathrm{T} 2}=\left(\frac{\rho_{20}}{\mathrm{~K}_{\mathrm{T}}}\right)\left(\frac{\mathrm{J}_{2}}{\mathrm{~J}_{1}}\right)\left(\mathrm{I}_{2} \mathrm{~W}_{2}\right)
$$

From here we find sizes ЛО:

$$
\begin{array}{r}
\mathrm{C}_{2}=\frac{\mathrm{K}_{\mathrm{T}} \tau_{\mathrm{T} 2}}{\mathrm{~K}_{32} \rho_{20} J_{2}^{2}}=\frac{\mathrm{K}_{\mathrm{T}}}{\mathrm{~K}_{32} \rho_{20} \mathrm{~J}_{2}^{2}}\left(\frac{\tau_{\text {ä }}}{\Delta_{\text {ok }}+\alpha \tau_{\text {aiii }}}\right) \\
\mathrm{h}_{2}=\frac{\rho_{20} \mathrm{~J}_{2}}{\mathrm{~K}_{\mathrm{T}} \tau_{\mathrm{T} 2}}\left(\mathrm{I}_{2} \mathrm{~W}_{2}\right)=\frac{\rho_{20} \mathrm{~J}_{2}}{\mathrm{~K}_{\mathrm{T}} \tau_{\text {âiì }}}\left(\Delta_{\text {ok }}+\alpha \tau_{\text {âiì }}\right)\left(\mathrm{I}_{2} \mathrm{~W}_{2}\right) \tag{5}
\end{array}
$$

As ampere coils are defined as:

$$
\mathrm{I}_{2} \mathrm{~W}_{2}=\mathrm{b}_{2} \mathrm{I}_{1} \mathrm{~W}_{1}=\mathrm{b}_{2} \sqrt{\frac{\mathrm{~K}_{\mathrm{u}} \mathrm{I}_{\mathrm{i}} \Delta \mathrm{U}}{\omega \mathrm{X}_{\mathrm{M}} \lambda}}
$$

That parity (5) can be written down as:

$$
\mathrm{h}_{2}=\frac{\mathrm{b}_{2} \rho_{20} \mathrm{~J}_{2}}{\mathrm{~K}_{\mathrm{T}} \tau_{\text {âï }}}\left(\Delta_{\mathrm{ok}}+\alpha \tau_{\text {âiii }}\right) \sqrt{\frac{\mathrm{K}_{\mathrm{u}} \mathrm{I}_{\mathrm{i}} \Delta \mathrm{U}}{\omega \mathrm{X}_{\mathrm{M}} \lambda}}
$$

Then the attitude of height of a winding to its thickness will be equal:

$$
\mathrm{n}_{\mathrm{e} 2}=\frac{\mathrm{h}_{2}}{\mathrm{C}_{2}}=\frac{\mathrm{J}_{2}^{3}}{\tau_{\mathrm{T} 2}^{2}}\left(\frac{\mathrm{~K}_{32} \rho_{20}}{\mathrm{~K}_{\mathrm{T}}}\right)^{2}\left(\mathrm{I}_{2} \mathrm{~W}_{2}\right)
$$

The received parities (4) - (5) show:

1. For essential reduction of thickness of winding $\mathrm{C}_{2}$ and a working backlash With it is necessary to increase density of current $\mathrm{J}_{2}$ and to reduce temperature of overheat $\tau_{\mathrm{adm}}$.
2. For reduction of height of a winding $h_{2}$ it is necessary to increase temperature of an overheat of winding $\tau_{\mathrm{adm}}$ and to reduce density of current $\mathrm{J}_{2}$ and ampere coils $\mathrm{I}_{2} \mathrm{~W}_{2}$.
3. The height of a winding $h_{2}$ is directly proportional ampere coils which in turn with increase in conductivity $\lambda$ and course $X_{M}$ decrease, and with increase in an increment of pressure $\Delta \mathrm{U}$ increase.
4. For elimination of harmony LW it is necessary to reduce density of current $\mathrm{J}_{2}$ and to increase temperature $\tau_{\mathrm{T}}$, that is first of all it is necessary to affect value of factor $\mathrm{n}_{\mathrm{e} 2}$.
Considering $\rho_{20}=1.7 \cdot 10^{-8}$ Ом м ; $\mathrm{K}_{32}=0.4-0.8 ; \mathrm{K}_{\mathrm{T}}=10-$ $13 ; \alpha_{\mathrm{M}}=0.0043 ; \mathrm{n}_{\mathrm{e} 2}=5-10 ; \tau_{2}=50-100^{\circ} \mathrm{C}$, it is easy to define factors $\mathrm{A}_{02} ; \Delta_{\mathrm{ok}}$ and thickness of winding $\mathrm{C}_{2}=\mathrm{C}^{\prime \prime}{ }_{2}$. But thus it is necessary to know value of density of current $\mathrm{J}_{2}$, satisfying as to a condition of a levitation, and set overheat $\tau_{2}$. Therefore all over again from conditions $\mathrm{C}_{2}^{\prime}=$ $\mathrm{C}^{\prime \prime}{ }_{2}$ we find density of current LW:

$$
\begin{equation*}
\mathrm{J}_{2}=\sqrt[3]{\left(\frac{\mathrm{K}_{32} \mathrm{n}_{\mathrm{e} 2} \mathrm{~A}_{\mathrm{o} 2}^{2}}{\mathrm{~b}_{2} \mathrm{~F}_{1}}\right) \frac{1}{\left(\alpha_{\mathrm{M}}+\frac{\Delta_{\mathrm{ok}}}{\tau_{2}}\right)^{2}}} \tag{6}
\end{equation*}
$$

Which satisfies formulas (3) and (6).
As $\mathrm{F}_{1}=\sqrt{\frac{2 \mathrm{P}_{\hat{\mathrm{a}}}}{\lambda}}=\sqrt{\frac{2 \mathrm{~K}_{\mathrm{u}} \mathrm{I}_{1}\left(\Delta \mathrm{U}_{1}\right)}{\omega \mathrm{X}_{\mathrm{M}} \lambda}}$,
That value (6) can be written down as:

$$
\begin{equation*}
\mathrm{J}_{2}=\dot{\mathrm{A}}_{0} \sqrt[3]{\frac{\mathrm{n}_{\mathrm{e} 2}}{\left(\alpha_{\mathrm{M}}+\frac{\Delta_{\mathrm{ok}}}{\tau_{2}}\right)^{2}} \sqrt{\frac{\lambda \mathrm{X}_{\mathrm{M}}}{\mathrm{I}_{1} \Delta \mathrm{U}_{1}}}} \tag{7}
\end{equation*}
$$

Where it is designated:

$$
\dot{A}_{0}=\sqrt[3]{\mathrm{A}_{\mathrm{o} 2}^{2} \frac{\mathrm{~K}_{32}}{\mathrm{~b}_{2}} \sqrt{\frac{\omega}{2 \mathrm{~K}_{\mathrm{u}}}}}
$$

The formula (7) allows to establish dependence:

$$
\mathrm{J}_{2}=\mathrm{f}\left(\mathrm{I}_{1} ; \mathrm{X}_{\mathrm{M}} ; \lambda ; \tau_{2} ; \mathrm{n}_{\mathrm{e} 2}\right)
$$



Fig. 4 Dependence of density of a current on a current at values of a course of mobile part

On fig. 4 dependence is resulted: $\mathrm{J}_{2}=\mathrm{f}\left(\mathrm{I}_{1}, \mathrm{X}_{\mathrm{M}}\right)$. The analysis of this dependence shows:

1. Density of a current in LW with increase in parameters $\mathrm{X}_{\mathrm{M}} ; \lambda ; \tau_{2} ; \mathrm{n}_{\mathrm{e} 2}$ increases.
2. With increase in rating values of stabilized current $\mathrm{I}_{1}$ the density of a current in $\mathrm{LW} \mathrm{J}_{2}$ falls.

## III. CONCLUSIONS

Analytical expressions of optimum values of parameters and the geometrical sizes of the precision stabilizer of an alternating current with use criterion of density of currents in windings Are received. The received expressions consider temperatures of an overheat of windings and course LW. Thus received following important conclusions:

1. The density of a current in LW turns out less, than density of a current in a winding of excitation (WE). For this reason the height appears less, than height LW.
2. The density of current LW is influenced essentially with values of height, a course and temperatures of overheat LW, and also specific conductivity of a working backlash. At increase in these parameters the density of a current in LW increases. However the increase in stabilizers of a current in WE leads to reduction of a current in LW. With reduction of thickness of a working backlash the density of a current in LW increases. Thus the increase in the attitude of height LW to its thickness is required.
3. At reduction of density of a current and increase in temperature of overheat LW the principle of harmony LW is carried out.

## REFERENCES

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